

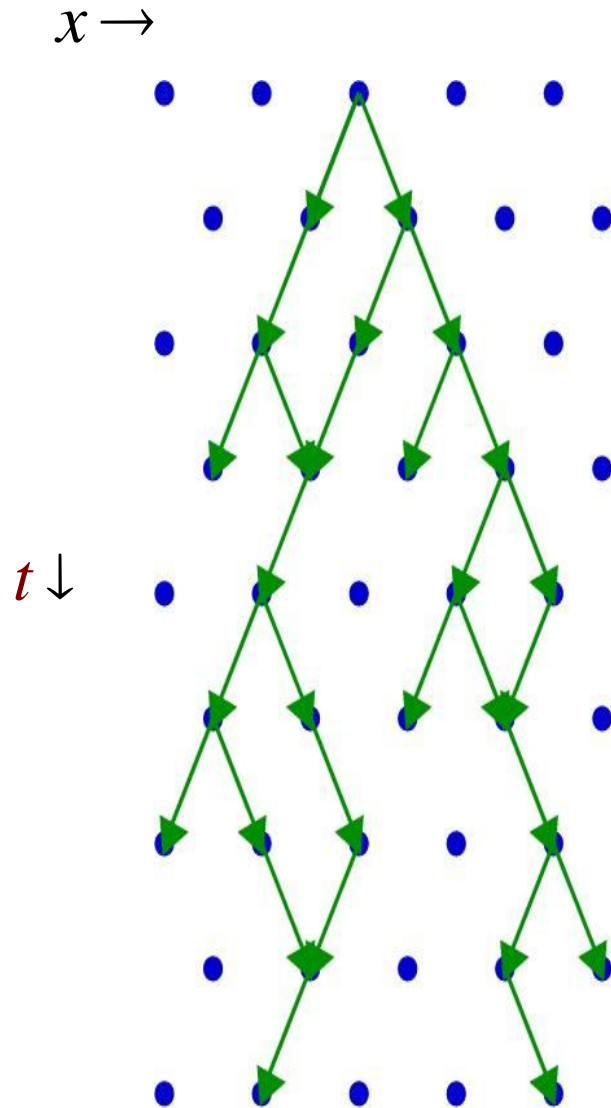
Critical behavior of a parity conserving cellular automaton with spatial disorder

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Motivation: finding nonequilibrium critical phase transition universality class candidate for experimental observation

Branching and annihilating random walks (BARW)

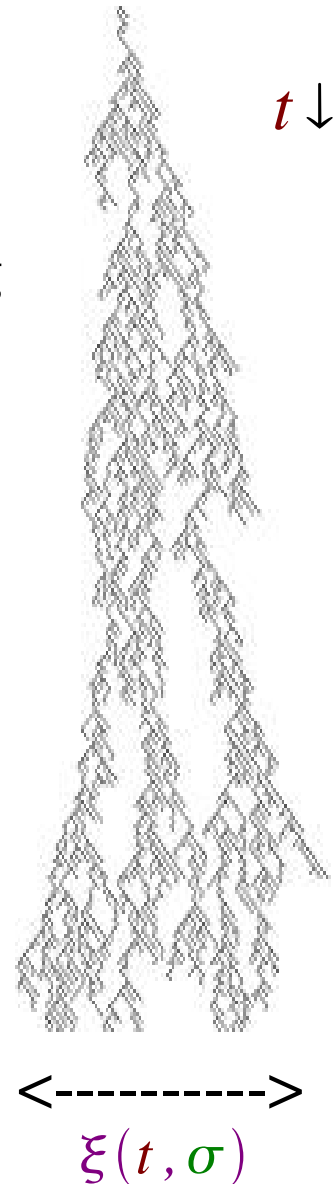


- $A \xrightarrow{\sigma} (n+1)A, \quad 2A \rightarrow A, \quad A \rightarrow 0$
 $\leftarrow 1+1 \text{ d realization} \rightarrow$
- For $n=1$ directed percolation (DP), contact process, epidemic spreading
 $\sim \text{BARW1} \sim \text{BARW}_{\text{odd}}$ (no $A \rightarrow 0$)
- Phase transition with DP class universal behavior:

$$\rho_{\infty} \propto (\sigma - \sigma_c)^{\beta} \quad \rho(t) \propto t^{-\alpha}$$

$$\xi_{\infty} \propto (\sigma - \sigma_c)^{-\nu} \quad R \propto \xi \propto t^{z/2}$$

- Well known exponents, field theory
- Robust universality class, still lacks experimental realization
 \leftarrow sensitivity to disorder

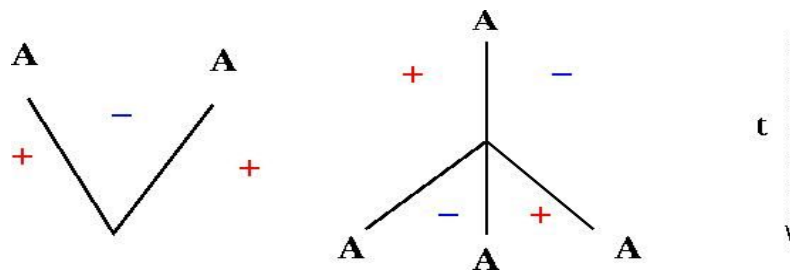


Annihilating random walks (ARW)

- $A + A \rightarrow A$ reaction ($\sim A + A \rightarrow 0$)
- Mean-field description: $d\rho(t, x)/dt = D\Delta\rho(t, x) - k\rho^2(t, x)$
asymptotic solution : $\rho(t) \sim 1/t$
- $d \leq 2$ spatial dimensions the fluctuations are relevant:
 $d = 1$: $\rho(t) \sim 1/\sqrt{t}$
 $d = 2$: $\rho(t) \sim \ln(t)/t$ “upper critical dimension”: $d_c = 2$
- Experimental verification in 1d TMMC:
R. Kroon, H. Fleurent, R. Sprik, PRE 47 (1993) 2462.

The parity conserving universality class (PC) BARW2, BARWe

- First seen in 1d stochastic CA by Grassberger 1984
- Nonequilibrium kinetic Ising model (NEKIM)
($T=0$ spin flip + $T>0$ spin flip) *N. Menyhárd, J. Phys. A 1994.*



Z_2 symmetry, domain walls \rightarrow particles: $AA \rightarrow 0$ $A \rightarrow 3A$
 parity conservation, non-DP class : $(AA \rightarrow 0, A \rightarrow 2A)$

- BARW2 or BARWe models Takayasu & Tretyakov, Jensen
- Field theoretical description Cardy & Tauber ('98) $d_c = 4/3$
NPRG by Canet et al. 2005

Knowledge about the effects of disorder (spatial quenched)

- Harris stability criterion: $d \nu > 2$
PC class: $d=1, \nu=1.857(1)$ -> **disorder relevant**
- Real space renormalization study of BARW2
(*Hooyberhs et al. 1999*): **No strong disorder fixed point**
- Inactive phase (ARW) :
 - Barriers and traps to diffusion -> subdiffusive behavior below a critical temperature, **slower nonuniversal power law**
Schütz & Mussawisade 1998 (exact calculation)
 - Random reaction probabilities -> **marginal perturbation**
Doussal & Monthous 1999 (real space renormalization)

NEKIMCA model algorithm

(Menyhárd & Ódor, JPA 1995)

- Diffusion: $\uparrow\uparrow\cdot\downarrow \xrightarrow{w_i} \uparrow\cdot\downarrow\downarrow$ ($\bullet = A$)
- Annihilation: $\uparrow\cdot\downarrow\cdot\uparrow \xrightarrow{w_o} \uparrow\uparrow\uparrow$
- Branching: $\uparrow\uparrow\cdot\downarrow\downarrow \xrightarrow{w_i^2} \uparrow\cdot\downarrow\cdot\uparrow\cdot\downarrow$

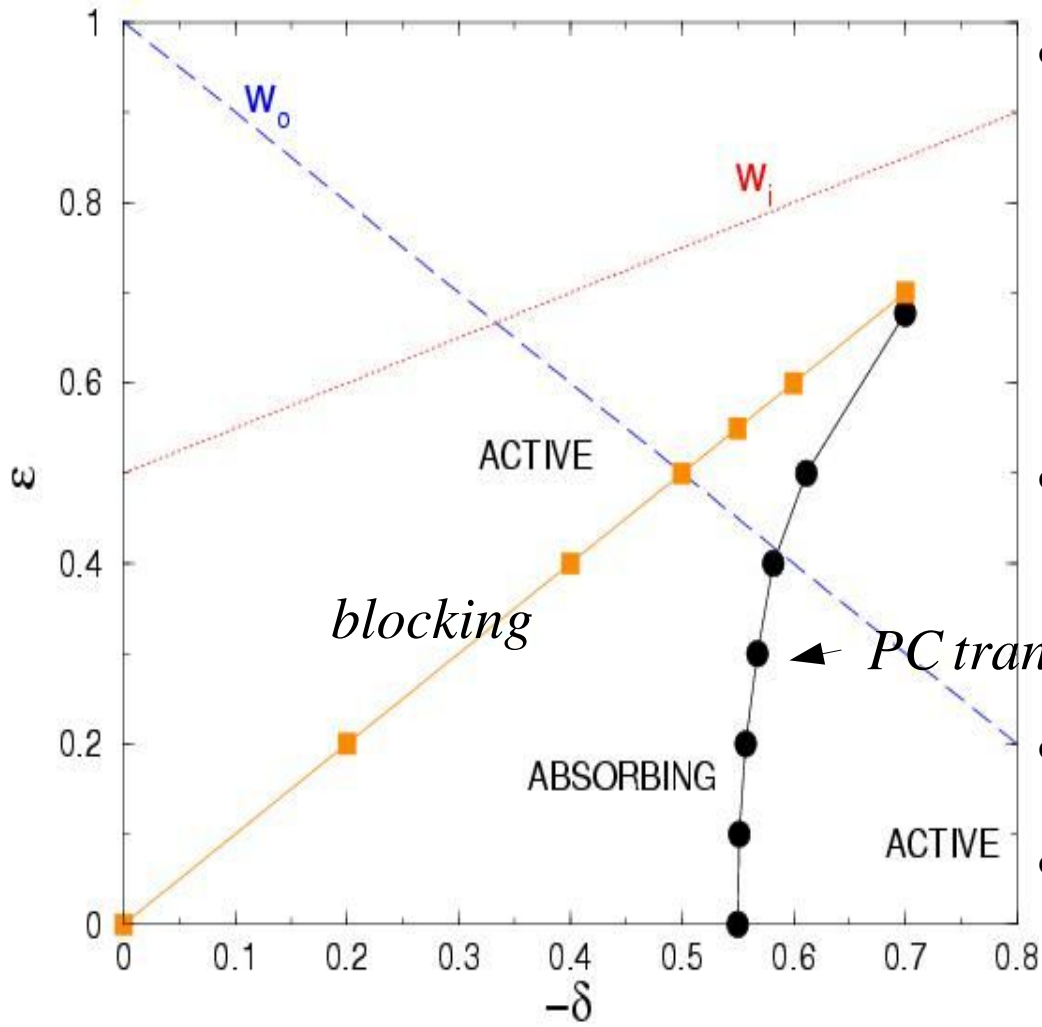
where:

$$w_i = (1 - \delta)/2 \qquad w_o = (1 + \delta)$$

The branching is the result of two overlapping diffusion steps in this synchronous SCA !

- Multispin (bit) coding, 32 or 64 samples evolve at once with the same random number sequence.

Results: Phase diagram



- Uniformly distributed disorder:

$$x(j) < q_i(j) = w_i + \chi(j) \quad x \in (0, 1)$$

$$y(j) < q_o(j) = w_o + \chi(j) \quad y \in (0, 1)$$

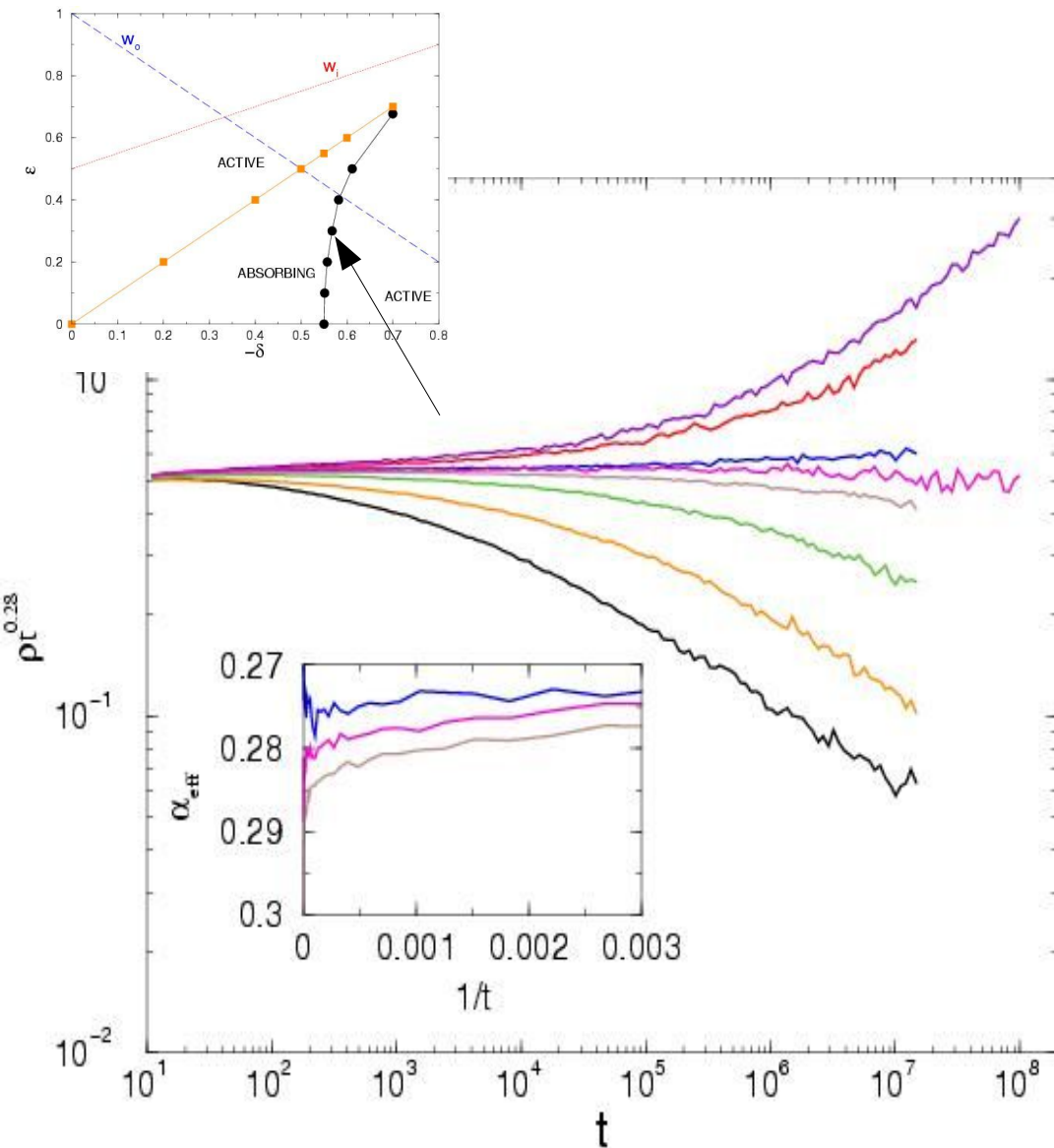
$$\chi \in (-\epsilon, \epsilon)$$

- Disorder increases the size of the absorbing phase

- Diffusion blocks where: $q_o(j) > 1$

- In odd parity blocks particles survive “blocking” transition to a new active phase

Scaling behavior on the PC class transition line I. (weak disorder $\epsilon < \sim 0.4$)



- Kink density decay from homogenous, random initial state, $L=10^5$, $t_{max}=10^8$ MCS

$$\rho \propto t^{-\alpha}$$

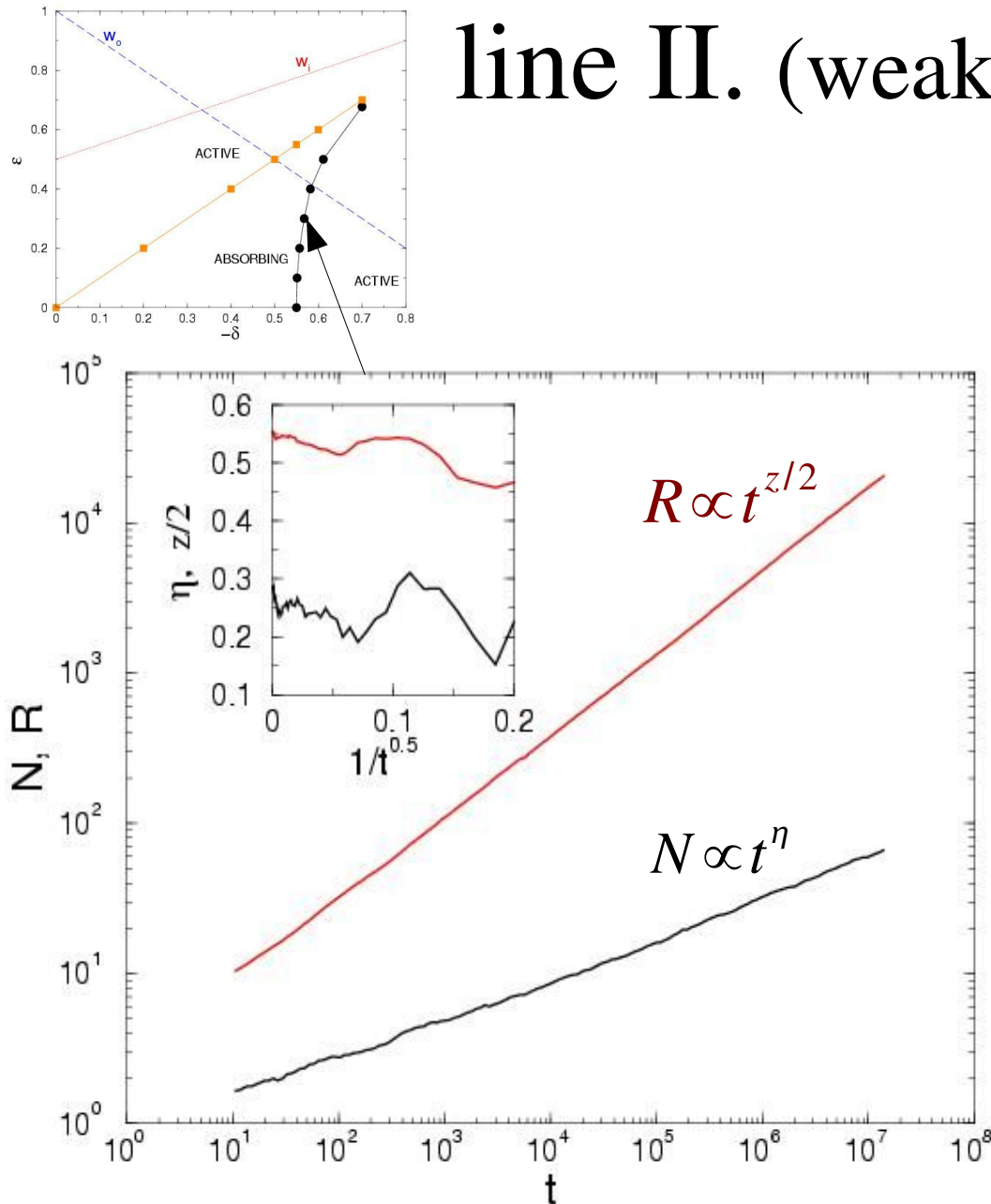
- PC class ($\alpha = 0.28(1)$) for weak disorder ($\epsilon = 0.0, 0.1, 0.2, 0.3, 0.4$)
- Effective exponent:

$$\alpha_{eff} = \frac{-\partial \ln(\rho)}{\partial \ln(t)}$$

$\epsilon = 0.3, -\delta = 0.571, 0.57, 0.568, 0.5675, 0.567, 0.565, 0.56, 0.55$

Scaling behavior on the PC-s class

line II. (weak disorder $\epsilon < \sim 0.4$)



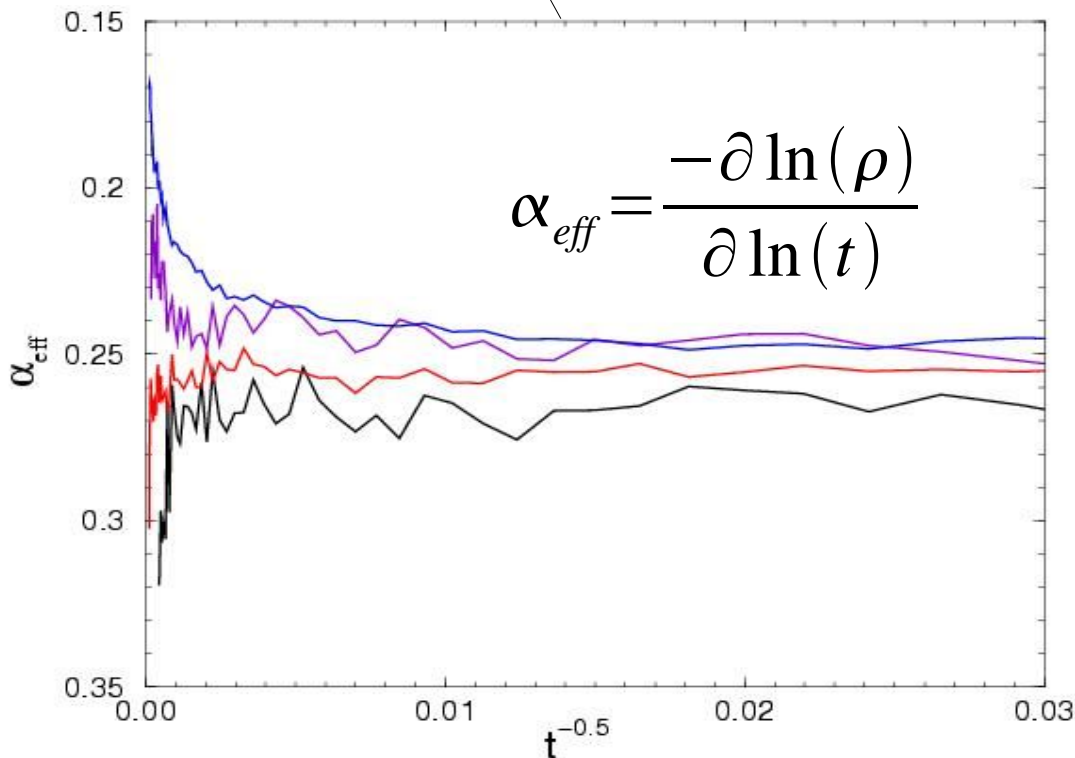
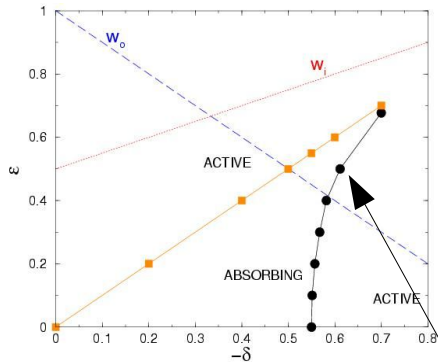
- Cluster evolution from a single kink, (odd sector),
 $L = 10^5$, $t_{\max} = 10^7$ MCS
- PC class:
 $\eta = 0.28(1)$, $z = 1.11(1)$
- Effective exponents:

$$\eta_{\text{eff}} = \frac{\partial \ln(N)}{\partial \ln(t)} \quad z/2_{\text{eff}} = \frac{\partial \ln(R)}{\partial \ln(t)}$$

$$\epsilon = 0.3, -\delta_c = 0.5679$$

Scaling behavior on the PC line III.

(strong disorder $\varepsilon > \sim 0.4$)



$\varepsilon = 0.5, -\delta = 0.613, 0.612, 0.611, 0.609$

- Kink density decay from homogenous, random initial state, $L = 10^5$, $t_{\max} = 10^8$ MCS
Blocking of annihilation \rightarrow continuously changing α ($\varepsilon = 0.5, 0.67$)
- Cluster exponents z, η decrease in this region.
- Rare region argument (ala Noest)

$$p_a \propto \exp(-ql_a), \quad (14)$$

$$\rho(t) \propto \int dl_a l_a p_a \exp[-t/\tau(l_a)], \quad (15)$$

$$\tau(l_a) \propto \exp(al_a) \quad (16)$$

$$\rho(t) \propto t^{-p_a/a}. \quad (17)$$

Scaling behavior below the PC line (active phase)

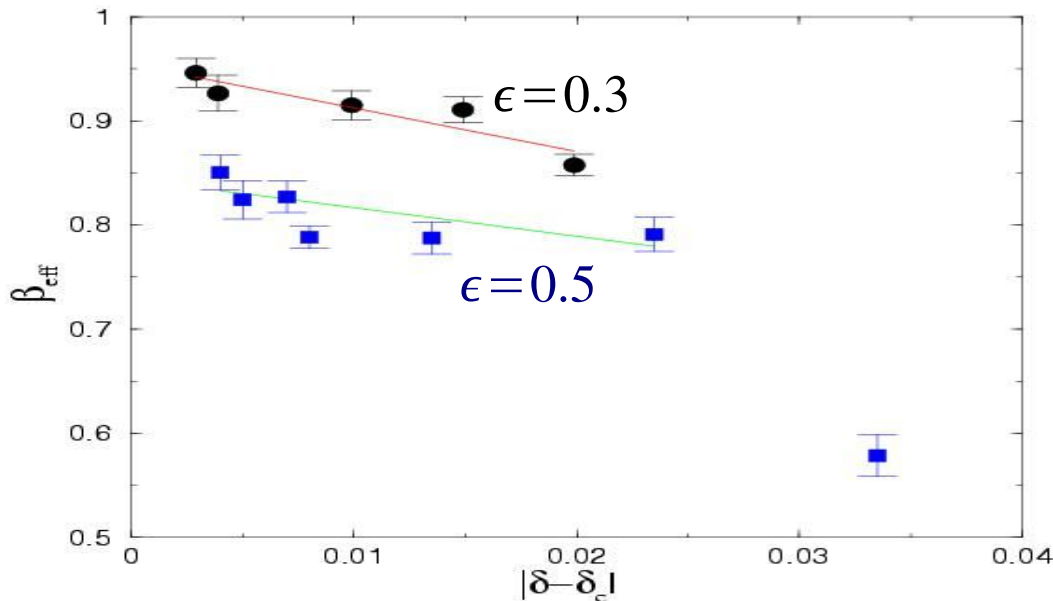
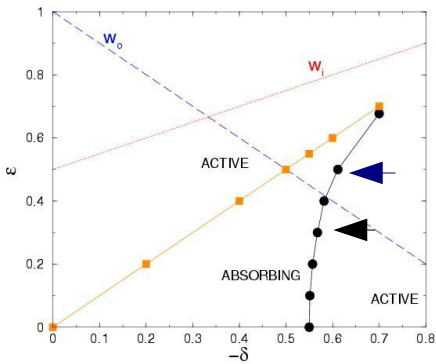


TABLE I. Numerical results for the disordered PC class transition line.

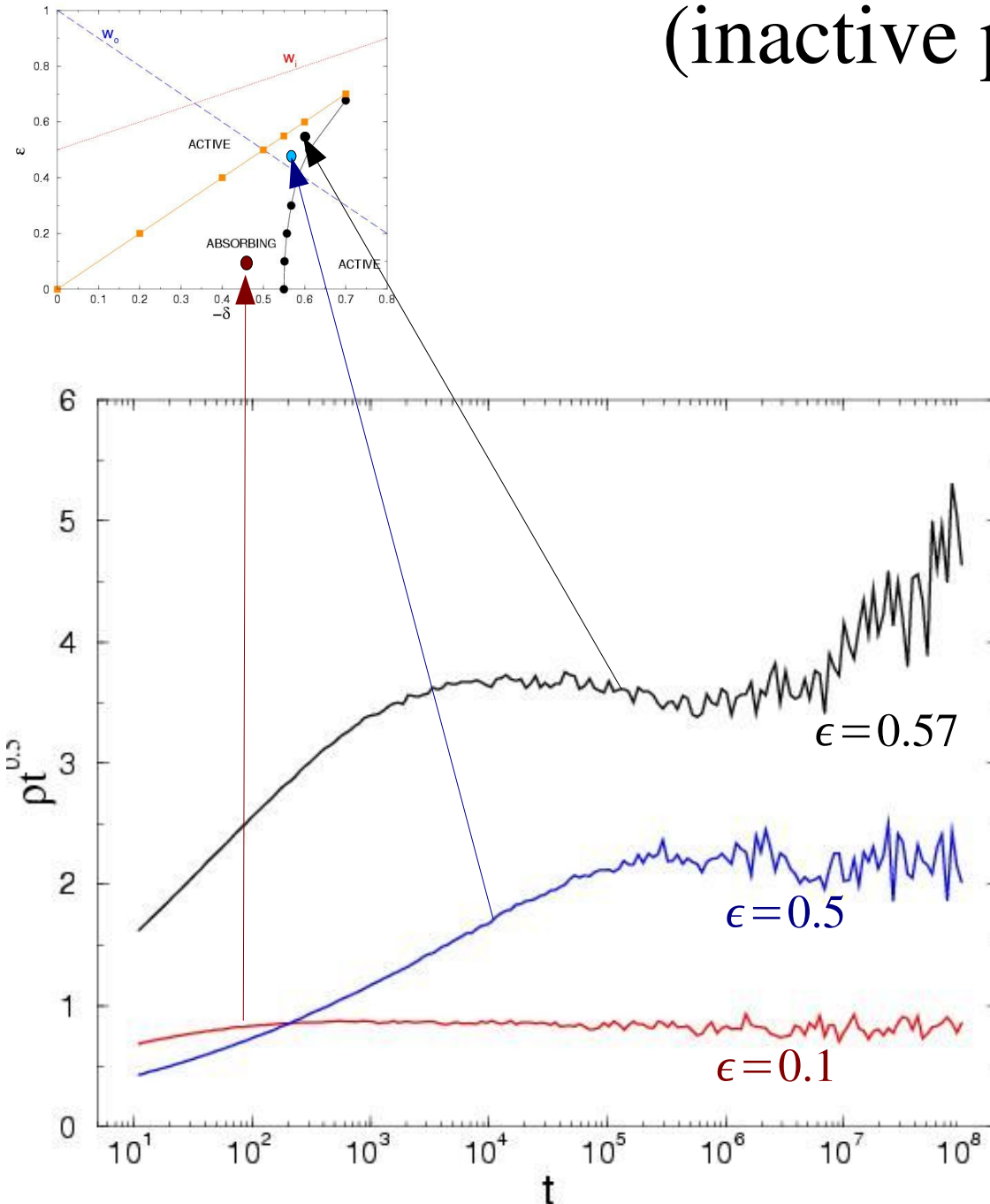
ϵ	$-\delta_c$	α	β	z	η
0.0	0.550(1)	0.280(6)			
0.1	0.5513(5)	0.280(6)			
0.2	0.557(1)	0.280(6)			
0.3	0.5676(1)	0.280(5)	0.95(1)	1.11(1)	0.285(5)
0.4	0.5849(1)	0.265(10)			
0.5	0.6115(4)	0.25(1)	0.84(2)	1.0(1)	0.252(6)
0.6767(2)	0.7	0.22(1)	0.80(2)		

- Steady state density effective exponent:

$$\beta_{eff} = \frac{-\partial \ln(\rho_\infty)}{\partial \ln(\epsilon)}$$

- PC class $\beta = 0.95(2)$
 $\epsilon < 0.4$ -re
- Continuously changing β , for strong disorder: $\epsilon > \sim 0.4$
Cluster exponents z, η decrease as well.

Scaling behavior above the PC line (inactive phase)

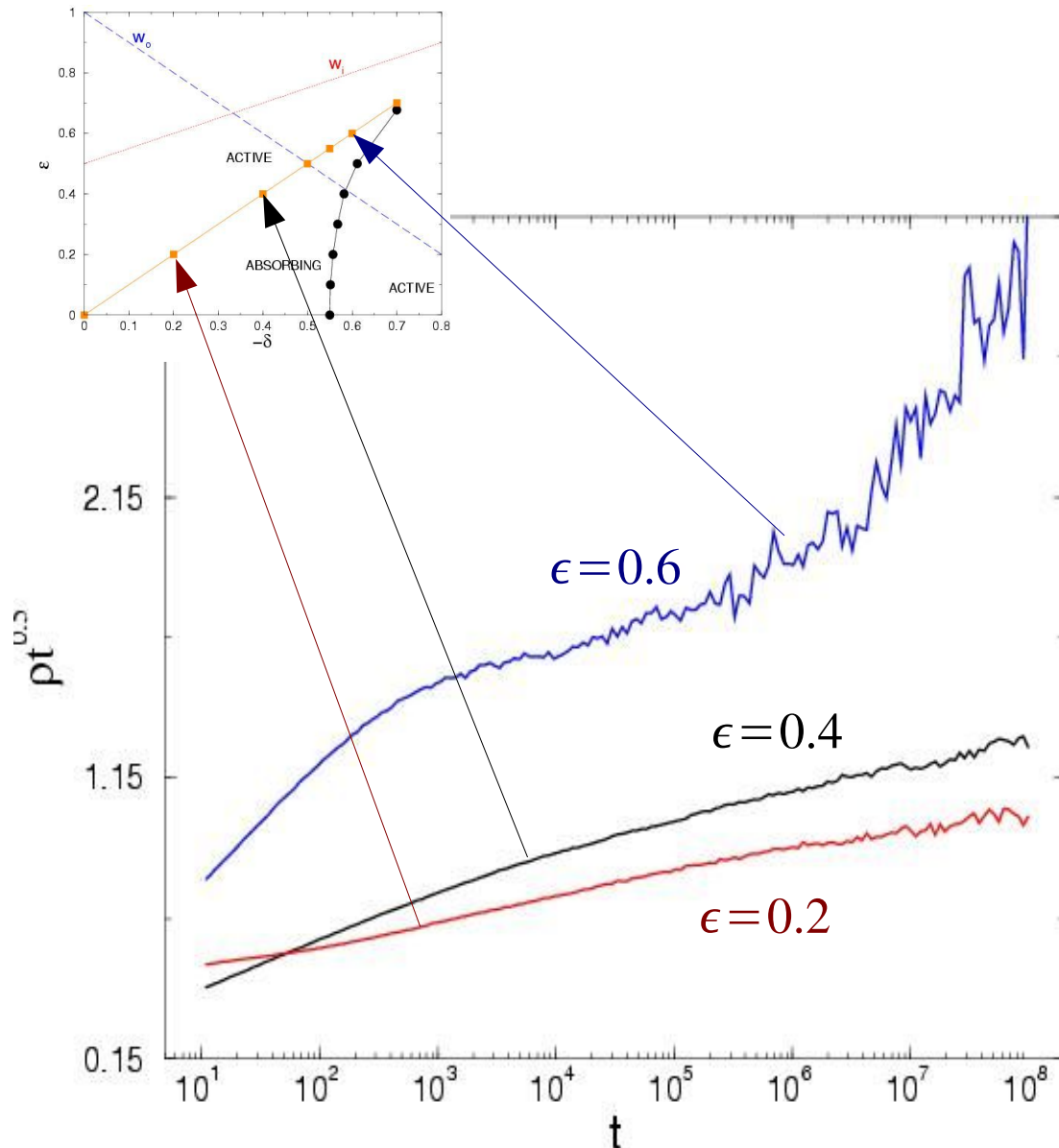


- Kink density decay from homogenous, random initial state, $L=10^5$, $t_{\max}=10^8$ MCS

$$\rho \propto t^{-\alpha}$$

- ARW class behavior $\alpha = 0.50(1)$ for small ϵ
- Increasing α just below the blocking transition line
- Agreement with Schütz's results

Scaling behavior on the blocking transition line (diffusion blocks) I.

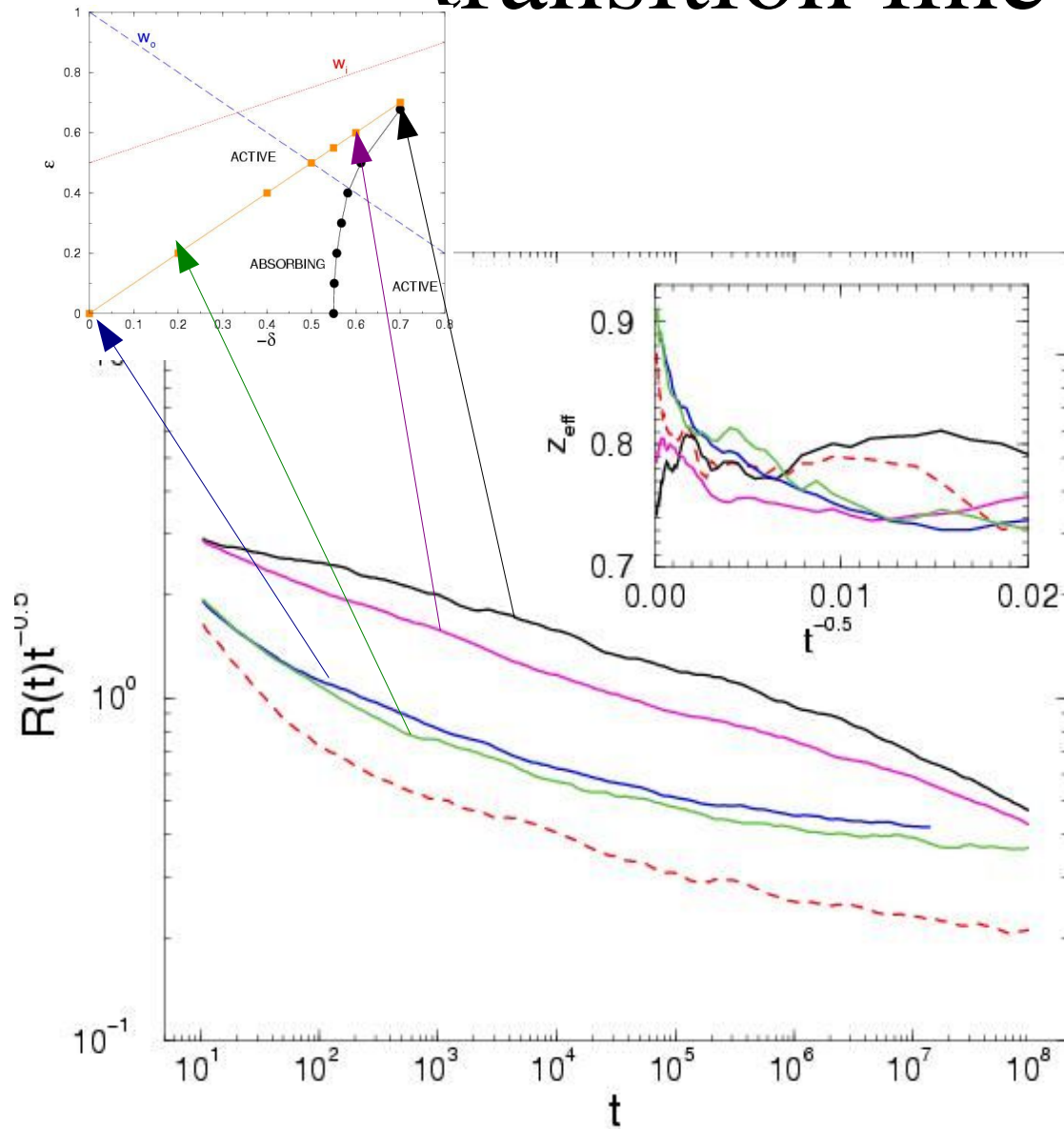


- Kink density decay from homogenous, random initial state, $L = 10^5$, $t_{\max} = 10^8$ MCS

$$\rho \propto t^{-\alpha}$$

- ARW + log. Corrections for weak disorder (without A-blocking)
- Increasing α (>0.5) if the annihilation is blocked as well.

Scaling behavior on the blocking transition line (clusters) II.



$\epsilon = 0.7, 0.6, 0.0, 0.2, 0.125$

- Cluster evolution from a single kink, (odd sector),
 $L=10^5$, $t_{\max}=10^8$ MCS

$$N \propto t^\eta \quad R \propto t^{z/2}$$

- ARW + corrections ($z=1$) for $\epsilon < 0.5$ (without annihilation blocking)
- Decreasing z if annihilation is blocked as well.
- $N(\infty)$ increases with ϵ , but $\eta=0$

Scaling above the blocking transition line (active fluctuating state)

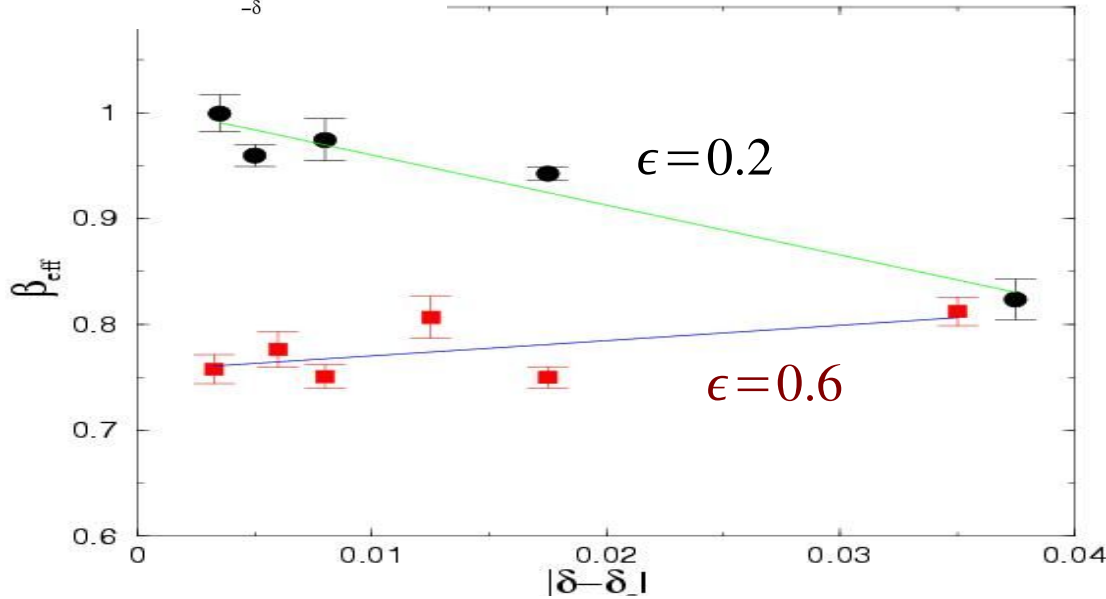
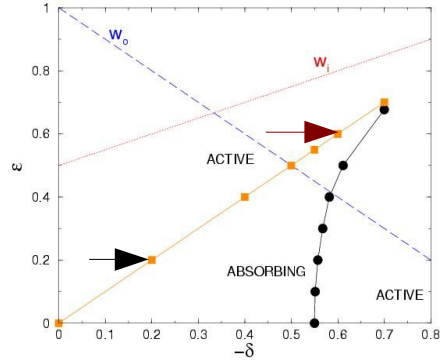


TABLE II. Numerical results for the freezing transition line.

ϵ	$-\delta_c$	α	β	z	$N(\infty)$
0.2	0.2	0.47(3)	1.00(5)	0.95(5)	1.031(1)
0.4	0.4	0.47(3)		0.95(5)	1.275(5)
0.5	0.5	0.43(1)		1.0(2)	1.85(5)
0.6	0.6	0.41(1)	0.75(3)	0.76(2)	3.75(5)
0.7	0.7	0.39(1)	0.68(4)	0.72(2)	14(1)

- Without annihilation blocking for $\epsilon < 0.5$: $\beta = 1$.

$$\rho \propto c_b \propto \langle l \rangle^{-1}, \quad (20)$$

For the exponential block size distribution (18) the average block size is

$$\langle l \rangle = [-\ln(p_{nw})]^{-1}, \quad (21)$$

with the no-wall probability

$$p_{nw} = 1 - p_w = 1 - (w_o + \epsilon - 1) = 2 - \epsilon - w_o. \quad (22)$$

The kink density is

$$\rho \propto -\ln(2 - \epsilon - w_o), \quad (23)$$

which for small wall probability p_w has the leading-order singularity

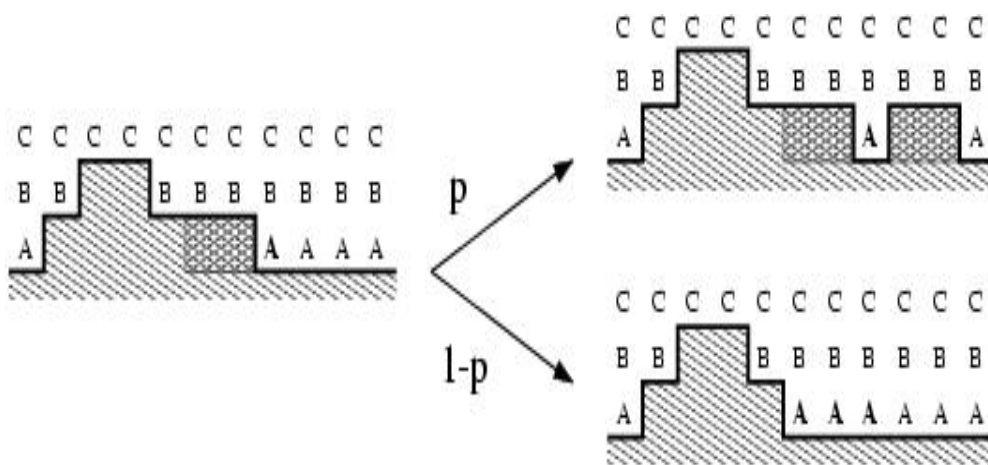
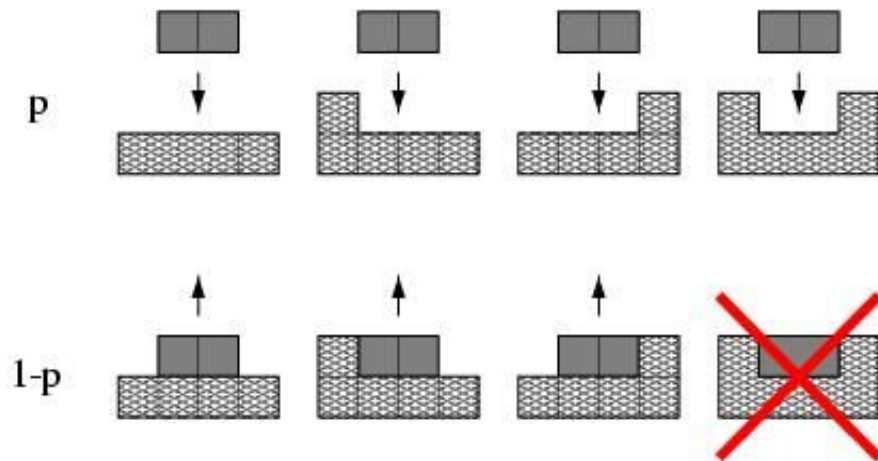
$$\rho \propto p_w. \quad (24)$$

- Decreasing β if the annihilation is blocked as well.

Summary

- A PC class for weak disorder -> **Experimental observation ?** For annihilation blocking continuously changing exponents
- $AA \rightarrow 0$ stable fixed point for weak disorder (inactive phase). For strong disorder the decay slows down (continuous exponents)
- “Blocking” transition appears for strong diffusion disorder
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- Acknowledgements: F. Iglói, G.M. Schütz, LCG, Clustergrid Hungary

Possible realization by a surface growth model



- Dimer growth model $\langle \leftrightarrow \rangle$
Unidirectionally coupled **parity conserving RD systems**:
 $A \rightarrow 3A$ $B \rightarrow 3B$ $C \rightarrow 3C \dots$
 $2A \rightarrow 0$ $2B \rightarrow 0$ $2C \rightarrow 0 \dots$
 $A \rightarrow A+B$ $B \rightarrow B+C$ $C \rightarrow C+D$
- Growth transition $\langle \leftrightarrow \rangle$
absorbing phase transition at each level with parity conserving class behavior